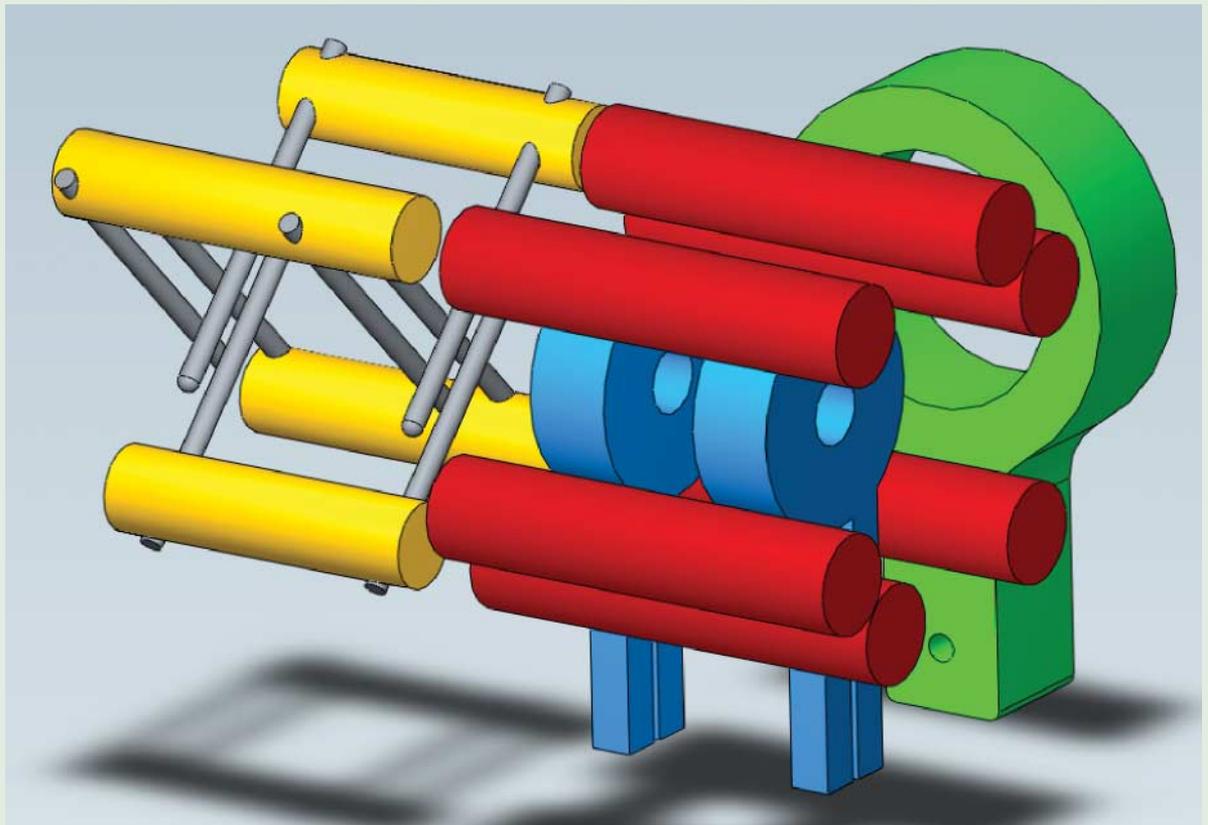


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## Magneto-electrostatic Trapping of Ground State OH Molecules

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We report magnetic confinement of neutral, ground state OH at a density of  $\sim 3 \times 10^3 \text{ cm}^{-3}$  and temperature of  $\sim 30 \text{ mK}$ . An adjustable electric field sufficiently large to polarize the OH is superimposed on the trap in various geometries, making an overall potential arising from both Zeeman and Stark effects. An effective molecular Hamiltonian is constructed, with Monte Carlo simulations accurately modeling the observed single-molecule dynamics in various trap configurations. Magnetic trapping of cold polar molecules under adjustable electric fields may enable study of low energy dipolar interactions.

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Cold and ultracold molecules are currently subject to intense research efforts as they promise to lead major advances in precision measurement, quantum control, and cold chemistry [1]. Much of this interest has focused on molecules possessing permanent electric dipole moments. The long range, anisotropic dipole-dipole interaction leads to novel collision [2,3] and reaction [4] dynamics, and may serve as the interaction between molecular qubits [5] in architectures for quantum information processing [6]. Quantum phase transitions can be explored with polar molecules [7]. Furthermore, polar molecules may possess extremely large internal electric fields (on the order of gigavolts per centimeter), yielding dramatic sensitivity enhancement to searches for an electron electric dipole moment [8,9]. Cold polar molecules have been produced via photoassociation of ultracold atomic species [10,11], Stark deceleration [12,13], and buffer gas cooling [14].

Polar molecules are readily trapped in inhomogeneous electric fields [15,16]; however, theory suggests that at high phase-space densities such traps invariably suffer from large inelastic collision losses [17]. One possible solution is to trap strong-field seeking states via time-dependent electric traps [18]. In contrast, we are interested in magnetically trapping these neutral polar molecules so as to have the freedom to apply external electric fields to control collision dynamics inside the trap. Highly vibrationally excited, triplet K<sup>+</sup>Rb molecules produced via photoassociation have been magnetically trapped at a temperature of  $\sim 300 \mu\text{K}$  and density of  $10^4 \text{ cm}^{-3}$  [19]. Buffer gas cooling has been used to load a magnetic trap with ground state CaH and NH molecules at temperatures of hundreds of millidegrees Kelvin and densities of  $\sim 10^8 \text{ cm}^{-3}$ , but in the presence of a cold He buffer gas [14,20]. We report here the first observation of magnetically trapped hydroxyl radical (OH) molecules and the lowest temperature (30 mK) yet achieved for a magnetically trapped polar molecule in its rovibronic ground state. Moreover, we perform the first study of trap dynamics based on a single molecule possessing large magnetic

and electric dipole moments in magnetic ( $B$ ) and electric ( $E$ ) fields that are inhomogeneous and anisotropic. Interesting polar molecule dynamics have been predicted at large field strengths [21]. Accurate correspondence between experimental data and simulations is achieved only by accounting for the molecular state mixing induced by the crossed fields. Large  $B$  fields may serve to suppress inelastic collision losses [22] while at the same time sympathetic cooling of molecules via cotrapped ultracold atoms may become feasible [23]. In addition, the magnetic quadrupole trap described here, in contrast to electrostatic traps, permits the application of arbitrary external  $E$  fields to a large class of polar molecules—a necessary step towards observing cold molecule dipole-dipole collisions [24].

To magnetically trap OH, we begin with a cold beam of ground state ( $^2\Pi_{3/2}$ ) OH produced via a pulsed electric discharge through 2% H<sub>2</sub>O vapor [25] seeded in 1.5 bar of Kr. A piezoelectric transducer valve [26] operating at 10 Hz provides a supersonic expansion of the OH/Kr mixture through a 1 mm nozzle. This results in a 490 m/s molecular packet with a 15% longitudinal ( $\hat{z}$ ) velocity spread. The transverse ( $\hat{\rho}$ ) velocity spread of the beam is limited to 2% by the 3 mm diameter molecular skimmer. Compared with a Xe expansion, the higher-velocity Kr expansion leads to less transverse beam loss within the current decelerator design. Thus, despite the fact that the Kr beam requires a higher phase angle for deceleration (smaller decelerator acceptance) it results in more decelerated molecules. After passing through the skimmer, the OH packet is spatially matched to the decelerator entrance via an electrostatic hexapole. Details regarding the OH Stark decelerator design and operating principle may be found in previous work [27,28]. We utilize the 142 stages of our Stark decelerator to slow a 130 mK packet (in the comoving frame) of OH to 20 m/s at a phase angle  $\phi_0 = 47.45^\circ$  for coupling into the magneto-electrostatic trap (MET). Decelerated OH packets are detected via laser-induced fluorescence (LIF) upon entering the MET, whose center lies 2.7 cm from the final pair of decelerator

rods. The 282 nm pulsed laser beam is oriented along the mutual longitudinal axis ( $\hat{z}$ ) of the decelerator and MET, and the resulting 313 nm LIF is collected in a solid angle of 0.016 sr.

The MET is illustrated in Fig. 1, and consists of two copper coils operated in an anti-Helmholtz configuration surrounded by an electric quadrupole. The center-to-center magnetic coil spacing is 1.5 cm, providing a field gradient of 6700 G/cm by passing 1500 A through the 4 turns of each water-cooled coil. A longitudinal trap depth of 400 mK is achieved, while the electrodes surrounding the coils allow the application of an  $E$  field to the trapped sample. The double- and single-rod electrodes that comprise the electric quadrupole allow the application of two distinct field configurations. A uniform bias electric field is created by grounding the single rods while oppositely charging the opposing double rods to  $\pm 3$  kV. Alternatively, we generate a quadrupole  $E$  field by charging neighboring rods with opposite polarities of  $\pm 9$  kV. Shielding by the MET coils reduces the interior  $E$  field by a factor of  $\sim 5$ . Figure 2 depicts the quadrupole electric field superimposed on the magnetic quadrupole field. OH has both an appreciable magnetic and electric dipole, and we can easily alter the geometry of the trapping potential by applying electric fields of variable magnitude and direction to the magnetically trapped sample.

The OH ground state is best described by Hund's case (a), in which the spin ( $|\Sigma| = \frac{1}{2}$ ) and orbital ( $|\Lambda| = 1$ ) angular momenta are strongly coupled to the molecular axis, yielding the total angular momentum projection  $|\Omega| = \frac{3}{2}$ . The  $|\Omega| = \frac{1}{2}$  state lies  $126 \text{ cm}^{-1}$  above this ground state. The nuclear spin of the hydrogen atom ( $I = \frac{1}{2}$ ) leads to a hyperfine splitting of the ground state into  $F =$

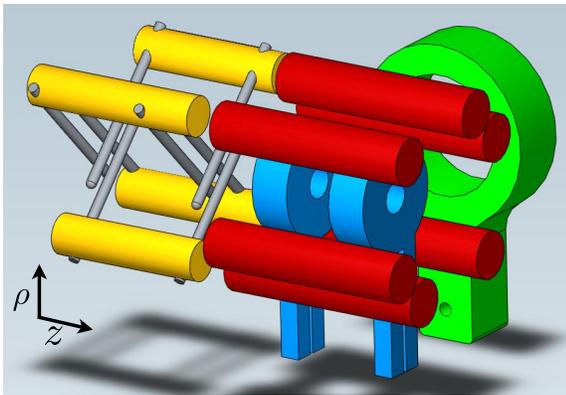


FIG. 1 (color online). Illustration of the magneto-electrostatic trap (MET) assembly. The terminus of the Stark decelerator, shown to the left, couples state-selected cold molecules into the magnetic quadrupole (a pair of rings in the middle). The double-electrode structure within the electric quadrupole (6 rods surrounding the magnetic quadrupole) allows for application of uniform electric fields within the magnetic trap. The ring at the back supports a 1 in. lens for collection of laser-induced fluorescence.

2, 1 components. In the presence of large external fields ( $E > 1 \text{ kV/cm}$  or  $B > 100 \text{ G}$ ), the electron angular momentum and nuclear spin decouple [5], and one considers the total angular momentum  $\mathbf{J}$ , which includes nuclear and electronic orbital angular momentum as well as electron spin. The projections  $m_J$  of  $\mathbf{J}$  are defined along the axis of the applied field. The ground state is further split into two opposite-parity states ( $f/e$ ) by the coupling of electronic orbital angular momentum to nuclear rotation. This small  $\Lambda$ -doublet splitting of 1.67 GHz, along with the 1.67 D electric dipole, is responsible for the rather large Stark shift experienced by OH molecules [29]. Figures 2(a) and 2(b) show the Zeeman and Stark effects of the OH ground state, respectively. External electric fields can thus dramatically alter the potential confining the magnetically trapped molecules. Figures 2(c)–2(e) illustrate the field geometries within the MET. In the plane perpendicular to  $\hat{z}$ , Fig. 2(c) shows the radial  $B$  field characteristic of a magnetic quadrupole, while Fig. 2(d) depicts the electric quadrupole field introduced by the electrode geometry. Figure 2(e) illustrates both  $B$  (small arrows) and  $E$  (large arrows) fields viewed from the side.

The top potential surface of Fig. 2(f) represents the energy shift of the decelerated  $|J = \frac{3}{2}, m_J = \frac{3}{2}\rangle$  state ( $E$

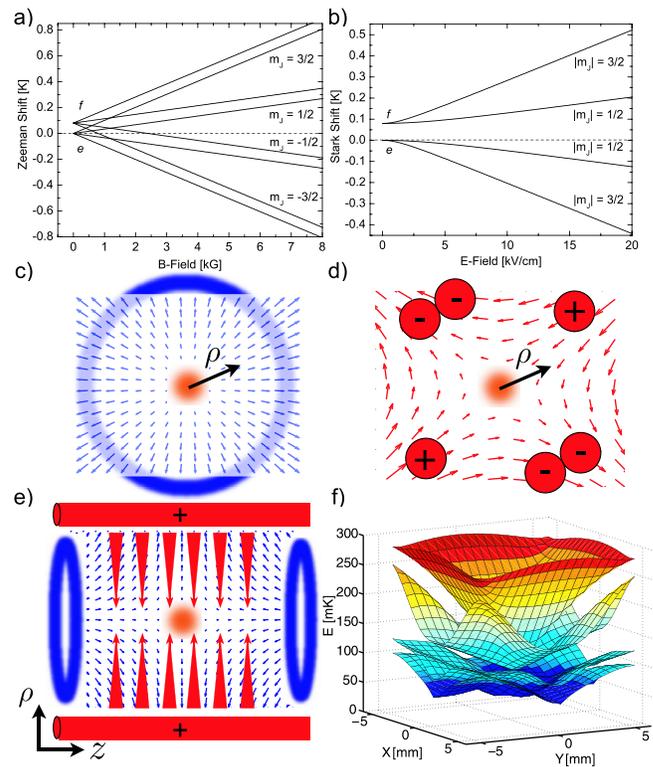


FIG. 2 (color online). (a) Zeeman and (b) Stark effects of the ground state structure of OH. (c) Magnetic and (d) electric quadrupole fields viewed from  $\hat{z}$ . (e) Side view of the MET configuration with quadrupole  $E$  and  $B$  fields. (f) Transverse adiabatic potential surfaces for various components of the OH ground state at the longitudinal MET center. The top surface depicts the decelerated and trapped  $|J = \frac{3}{2}, m_J = \frac{3}{2}\rangle$  state.

fields couple  $e$  and  $f$ , therefore they are no longer good quantum numbers) within the combined quadrupole  $E$  and  $B$  fields present in the MET. The transverse profiles of the four trapping adiabatic potentials are depicted at the longitudinal trap center, where  $\vec{B} \cdot \hat{z} = 0$ . The remaining four surfaces (not shown) are simply inverted relative to those displayed, and are therefore antitrapping. The three-dimensional MET potentials are calculated by diagonalizing an effective Hamiltonian that includes both Stark and Zeeman terms. Coupling between the ground  $|J = \frac{3}{2}, \Omega = \frac{3}{2}\rangle$  and excited  $|\frac{5}{2}, \frac{3}{2}\rangle$ ,  $|\frac{1}{2}, \frac{1}{2}\rangle$ ,  $|\frac{3}{2}, \frac{1}{2}\rangle$  states is included, for a total of 64 hyperfine components. The  $|J = \frac{3}{2}, \Omega = \frac{3}{2}\rangle$ ,  $|\frac{3}{2}, \frac{1}{2}\rangle$  coupling is the most critical, increasing the trap depth of the top surface of Fig. 2(f) by 11%. Addition of the remaining states modifies the potential by  $<0.1\%$ . The square shape of the top surface is a result of the varying angle between  $E$  and  $B$  in the MET. The potentials from each field directly sum where the fields are collinear, but the trap shallows from this maximum value as the angle between  $E$  and  $B$  increases.

The Stark decelerator slows only those OH molecules in the electrically weak-field seeking (EWFS) component of the ground state, i.e., the upper  $\Lambda$  doublet. Of these EWFS molecules, only those with components  $J = \frac{3}{2}$ ,  $m_J = \pm \frac{3}{2}$  are synchronously slowed to 20 m/s. Because the positive  $m_J$  state is magnetically weak-field seeking, only 50% of the decelerated molecular packet is trapable by the MET. The molecules experience a  $B$  field  $>100$  G at the last deceleration stage emanating from the back coil, preserving the quantization of the decelerated state.

To trap the incoming OH packet, the coil closer to the decelerator (front coil) is grounded while the coil further from the decelerator (back coil) operates at 2000 A for the duration of the 3.7 ms deceleration sequence. The effective magnetic dipole of OH in its  $J = m_J = \frac{3}{2}$  state is  $1.2\mu_B \approx 0.81$  K/T, where  $\mu_B$  is the Bohr magneton, allowing the back coil to stop molecules with velocity  $\leq 22$  m/s. The front coil is turned on after the 20 m/s molecules are stopped by the back coil, which occurs 2.65 ms after exiting the final deceleration stage. The trap switching is synchronous to the 20 m/s molecules since OH possessing larger longitudinal velocities escape the quadrupole field. For technical reasons, the two magnet coils are connected in series, causing the current in both coils to decrease from 2000 to 1500 A over  $\sim 800 \mu\text{s}$  once the front coil is switched into the circuit.

Figure 3 displays typical time-of-flight data resulting from the trapping sequence along with results from Monte Carlo simulation of the trap dynamics. The large initial peak is the decelerated OH packet traversing the center of the trap. All subsequent peaks are the result of oscillation in the  $\hat{\rho}$  dimension, with a frequency of  $\sim 650$  Hz. Fluorescence is gathered in  $\hat{\rho}$ , and oscillations in  $\hat{z}$  are not detectable since our detection region spans the entire visible trap length. The magnetic quadrupole trap is unable to confine all the molecules stopped by the back coil

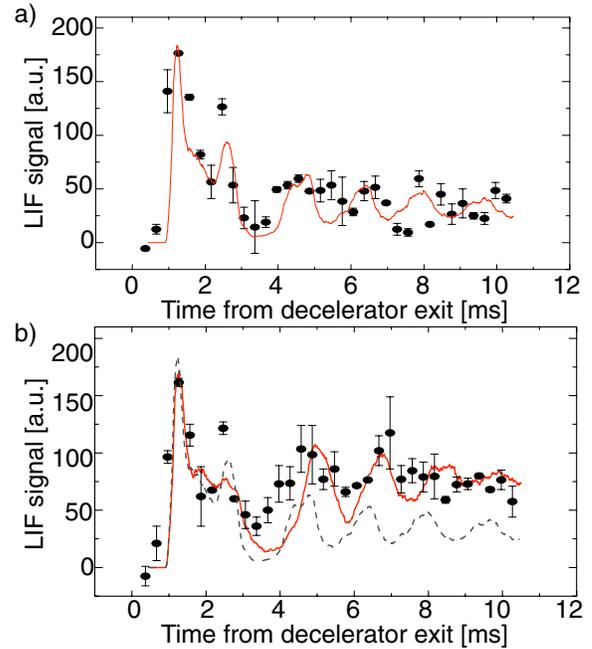


FIG. 3 (color online). Time-of-flight data and Monte Carlo simulation results (solid red line) for two different MET configurations. The magnetic field switches from the stopping configuration to quadrupole trapping at  $t = 2.65$  ms. (a) Magnetic trap only. (b) Stopping and trapping in the presence of combined electric and magnetic quadrupole fields. Note the larger steady-state trap population and damped oscillations compared to the zero  $E$ -field case (dashed gray line).

alone due to the 25% current drop, as is clear from the difference in signal height at  $\sim 2.5$  ms versus the steady-state level at, e.g., 10 ms. Note the larger steady-state population of Fig. 3(b) from the addition of a confining transverse electric quadrupole. Also visible is the shorter coherence time compared to Fig. 3(a), which results from trap distortion visible in Fig. 2(f). An estimate of the trap density, as measured by LIF, gives  $\sim 3 \times 10^3 \text{ cm}^{-3}$ . By comparison, the density of a packet “bunched” at a phase angle of  $\phi_0 = 0^\circ$  is  $10^7 \text{ cm}^{-3}$ , while the 20 m/s packet is  $10^4 \text{ cm}^{-3}$  due to its spreading during free-flight to the trap center. Previous work with electrostatic traps yielded  $\sim 10^3$  higher densities of trapped OH [15]. The trap density can be further increased by decreasing the distance between decelerator and MET or by applying larger voltages to the electric quadrupole. Monte Carlo analysis of the velocity distribution of the trapped OH yields temperatures of  $\sim 30$  mK.

Higher decelerator phase angles should, in principle, produce a slower OH packet. However, below 20 m/s the OH population drops sharply and no further deceleration is observed. We attribute this effect to a reduction in the effective aperture seen by slow molecules as they traverse the final deceleration stage. The OH closest to the final decelerator rods see a larger longitudinal potential than those on axis and, at sufficient phase angles, are stopped or

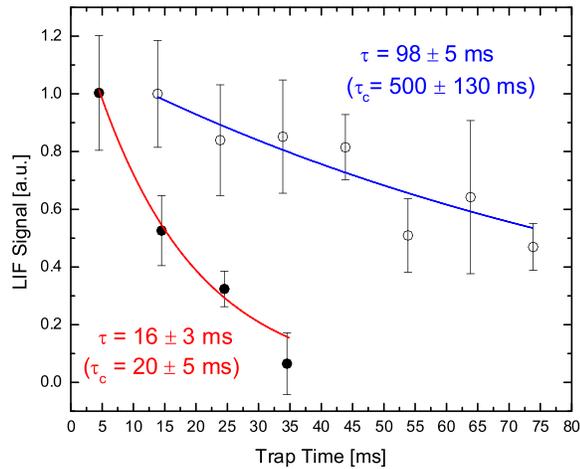


FIG. 4 (color online). Measured ( $\tau$ ) and deconvolved collisional ( $\tau_c$ ) lifetimes of magnetically trapped OH at a background pressure of  $1 \times 10^{-6}$  Torr (●) and  $4 \times 10^{-8}$  Torr (○).

reflected and hence cannot be observed in the MET region. Numerical models elucidate this effect, which may be mitigated by new decelerator designs [30].

The 10 Hz repetition rate of our Stark decelerator currently limits our trap interrogation time to 100 ms. Maintaining continuous currents of 1500 A presents technical challenges due to heating in both the switching transistors and cabling, which require water cooling. Figure 4 shows electric field-free OH trap lifetimes measured in two different background pressure regimes:  $1 \times 10^{-6}$  and  $4 \times 10^{-8}$  Torr, both dominated by  $N_2$  at 295 K. Nonideal power supply control causes the trap current to decrease from 1500 to 1160 A at a rate of 2000 A/s. This lowers the trap depth, resulting in a loss of  $12 \pm 2 \text{ s}^{-1}$ . After accounting for this loss mechanism, we find background collision-limited lifetimes of  $20 \pm 5$  and  $500 \pm 130$  ms for the two pressure regimes, respectively. These lifetimes determine an OH- $N_2$  collision cross section of  $500 \pm 100 \text{ \AA}^2$ . From this clear dependence of trap lifetime on background gas pressure, we envision measuring dipole-dipole collision cross sections by using another beam of polar molecules colliding under the  $E$  field of the MET.

We have magnetically confined ground state hydroxyl radicals in a MET that allows for the application of a variable  $E$  field to the trapped polar molecules. Diagonalization of an effective Hamiltonian modeling the energy shift of OH in combined  $E$  and  $B$  fields was necessary for accurate Monte Carlo simulation of the trap dynamics. Future experiments could quantify temperature dependent nonadiabatic transition rates between the different surfaces depicted in Fig. 2(f) akin to Majorana transitions observed in atomic magnetic quadrupole traps. Furthermore, higher densities within the MET will facilitate the study of cold OH-OH collisions in variable electric fields. As rather

modest electric fields are expected to modify collision cross sections between polar molecules by as much as  $10^3$  [17], crossed-beam experiments exploiting the MET described here will benefit from both low center-of-mass collision energies and tunable  $E$  fields in the interaction region.

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- [1] J. Doyle *et al.*, Eur. Phys. J. D **31**, 149 (2004).
- [2] A. V. Avdeenkov, D. C. E. Bortolotti, and J. L. Bohn, Phys. Rev. A **69**, 012710 (2004).
- [3] J. J. Gilijamse *et al.*, Science **313**, 1617 (2006).
- [4] E. R. Hudson *et al.*, Phys. Rev. A **73**, 063404 (2006).
- [5] B. L. Lev *et al.*, Phys. Rev. A **74**, 061402(R) (2006).
- [6] D. DeMille, Phys. Rev. Lett. **88**, 067901 (2002).
- [7] A. Micheli, G. Brennen, and P. Zoller, Nature Phys. **2**, 341 (2006).
- [8] J. J. Hudson *et al.*, Phys. Rev. Lett. **89**, 023003 (2002).
- [9] M. G. Kozlov and D. DeMille, Phys. Rev. Lett. **89**, 133001 (2002).
- [10] K. M. Jones *et al.*, Rev. Mod. Phys. **78**, 483 (2006).
- [11] J. Sage *et al.*, Phys. Rev. Lett. **94**, 203001 (2005).
- [12] H. L. Bethlem, G. Berden, and G. Meijer, Phys. Rev. Lett. **83**, 1558 (1999).
- [13] J. Bochinski *et al.*, Phys. Rev. Lett. **91**, 243001 (2003).
- [14] J. Weinstein *et al.*, Nature (London) **395**, 148 (1998).
- [15] S. Y. T. van de Meerakker *et al.*, Phys. Rev. Lett. **94**, 023004 (2005).
- [16] T. Rieger *et al.*, Phys. Rev. Lett. **95**, 173002 (2005).
- [17] A. Avdeenkov and J. Bohn, Phys. Rev. A **66**, 052718 (2002).
- [18] J. van Veldhoven, H. L. Bethlem, and G. Meijer, Phys. Rev. Lett. **94**, 083001 (2005).
- [19] D. Wang *et al.*, Phys. Rev. Lett. **93**, 243005 (2004).
- [20] W. C. Campbell *et al.*, arXiv:physics/0702071 [Phys. Rev. Lett. (to be published)].
- [21] T. V. Tscherebul and R. Krems, Phys. Rev. Lett. **97**, 083201 (2006).
- [22] C. Ticknor and J. L. Bohn, Phys. Rev. A **71**, 022709 (2005).
- [23] M. Lara *et al.*, Phys. Rev. A **75**, 012704 (2007); P. Soldan and J. M. Hutson, Phys. Rev. Lett. **92**, 163202 (2004).
- [24] R. V. Krems, Phys. Rev. Lett. **96**, 123202 (2006).
- [25] H. J. Lewandowski *et al.*, Chem. Phys. Lett. **395**, 53 (2004).
- [26] D. Proch and T. Trickl, Rev. Sci. Instrum. **60**, 713 (1989).
- [27] J. Bochinski *et al.*, Phys. Rev. A **70**, 043410 (2004).
- [28] E. R. Hudson *et al.*, Eur. Phys. J. D **31**, 351 (2004).
- [29] E. R. Hudson *et al.*, Phys. Rev. Lett. **96**, 143004 (2006).
- [30] B. Sawyer *et al.*, arXiv:physics/07053442.